## Short Note

# A proof that a discrete delta function is second-order accurate 

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#### Abstract

It is proved that a discrete delta function introduced by Smereka [P. Smereka, The numerical approximation of a delta function with application to level set methods, J. Comput. Phys. 211 (2006) 77-90] gives a second-order accurate quadrature rule for surface integrals using values on a regular background grid. The delta function is found using a technique of Mayo [A. Mayo, The fast solution of Poisson's and the biharmonic equations on irregular regions, SIAM J. Numer. Anal. 21 (1984) 285-299]. It can be expressed naturally using a level set function. © 2007 Elsevier Inc. All rights reserved.


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There is considerable interest in designing accurate discrete delta functions for surfaces in a domain covered by a rectangular grid. They can provide quadrature rules for surface integrals using values at regular grid points [2,10-12]. Such a rule is especially useful when the surface is represented by a level set function. In [10] Smereka constructed a discrete delta function as the truncation error in applying the discrete Laplacian to a "Green's function" for the exact delta function on the surface. To find the truncation error, he used the technique of Mayo $[7,8]$ for solving differential equations with interfacial conditions, in which jump conditions are built into the difference operators on a regular grid. (The immersed interface method [3,5], the EJIIM [13,9] and the ghost fluid method [6] are related to Mayo's technique.) Smereka also showed how to express this delta function in terms of a level set function. He conjectured that the resulting quadrature rule for surface integrals is second-order accurate and verified the accuracy in numerical examples. In this note we give a simple proof of this fact.

Suppose $\Gamma$ is a closed curve in $\mathbb{R}^{2}$ or a closed surface in $\mathbb{R}^{3}$, bounding a set which is contained in a rectangular domain $\Omega$. The problem is to design a weight function $w^{\text {h }}$ at grid points on a square grid $\Omega_{h}$, concentrated near $\Gamma$, so that, for any smooth function $f$ defined near the curve $\Gamma$ in $\mathbb{R}^{2}$,

[^0]\[

$$
\begin{equation*}
\int_{\Gamma} f(x) \mathrm{d} s(x)=\sum_{i h \in \Omega_{h}} f(i h) w^{h}(i h) h^{2}+\mathrm{O}\left(h^{2}\right) \tag{1}
\end{equation*}
$$

\]

or near the surface $\Gamma$ in $\mathbb{R}^{3}$,

$$
\begin{equation*}
\int_{\Gamma} f(x) \mathrm{d} S(x)=\sum_{i h \in \Omega_{h}} f(i h) w^{h}(i h) h^{3}+\mathrm{O}\left(h^{2}\right) . \tag{2}
\end{equation*}
$$

Arclength and surface area are special cases. Smereka's $w^{\mathrm{h}}$ has support on the grid points $i h$ within distance $h$ of $\Gamma$, i.e., $w^{h}(i h)=0$ at other points. We will prove that (2) holds, with $w^{\mathrm{h}}$ as in [10], assuming $\Gamma$ is a smooth surface in $\mathbb{R}^{3}$. The case of a curve in $\mathbb{R}^{2}$ is entirely similar.

Smereka's procedure is as follows: Let $\delta_{\Gamma}$ be the distribution, or generalized function, restricting to $\Gamma$; that is, for smooth $f$ on $\Omega$,

$$
\begin{equation*}
\int_{\Omega} f \delta_{\Gamma} \mathrm{d} x=\int_{\Gamma} f \mathrm{~d} S . \tag{3}
\end{equation*}
$$

Let $g$ be the solution of

$$
\begin{equation*}
\Delta g=\delta_{\Gamma} \quad \text { in } \Omega, \quad g=0 \quad \text { on } \quad \partial \Omega . \tag{4}
\end{equation*}
$$

Assuming $\Gamma$ is smooth, $g$ is piecewise smooth, i.e., smooth and harmonic on each region bounded by $\Gamma$, with the jump conditions

$$
\begin{equation*}
[g]=0, \quad\left[\partial_{n} g\right]=1 \text { on } \Gamma, \tag{5}
\end{equation*}
$$

where $\partial_{n}$ is the normal derivative on $\Gamma$. In fact $g$ can be thought of as a single layer potential on $\Gamma$. Now let $\Delta_{h}$ be the usual second-order discrete Laplacian on $\Omega_{h}$ and let $\tau^{h}$ be the truncation error

$$
\begin{equation*}
\Delta_{h} g=\tau^{h} \quad \text { on } \Omega_{h} \tag{6}
\end{equation*}
$$

Smereka constructs the weights $w^{\mathrm{h}}$ from expressions for $\tau^{h}$, using Mayo's technique [7,8]. At a regular grid point $i h \in \Omega_{h}$, for which the stencil of $\Delta_{h}$ does not cross $\Gamma, \tau^{h}(i h)=\mathrm{O}\left(h^{2}\right)$ as usual. At an irregular grid point, $\tau^{h}$ is larger. It can be found to $\mathrm{O}(h)$ using the jumps in first and second derivatives of $g$; see (30) in [10]. These in turn can be expressed in derivatives of the normal and tangent vectors to $\Gamma$. (See (41), (47) in [10] for $\mathbb{R}^{2}$ and Section 7.2 for $\mathbb{R}^{3}$.) Thus $\tau^{h}$ has the form

$$
\begin{equation*}
\Delta_{h} g=\tau^{h}=w^{h}+\mathrm{O}_{\Gamma}(h)+\mathrm{O}\left(h^{2}\right) \quad \text { on } \Omega_{h}, \tag{7}
\end{equation*}
$$

where $w^{\mathrm{h}}$ is known analytically and $w^{h}$ and $\mathrm{O}_{\Gamma}(h)$ are nonzero only at the irregular points. The errors are uniform. Smereka shows how to write $w^{\mathrm{h}}$ in terms of a level set function; see (45) and Section 7 in [10].

To prove that (2) is valid, we may assume $f$ is nonzero only in a neighborhood of $\Gamma$, as well as smooth. We begin by writing

$$
\begin{equation*}
\int_{\Gamma} f \mathrm{~d} S=\int_{\Omega} f \delta_{\Gamma} \mathrm{d} x=\int_{\Omega} f \Delta g \mathrm{~d} x=\int_{\Omega} g \Delta f \mathrm{~d} x . \tag{8}
\end{equation*}
$$

(This could be rewritten in an equivalent way using the jump conditions (5) rather than $\delta_{\Gamma}$ ).
Next we replace the last integral by a sum over grid points. We check that

$$
\begin{equation*}
\int_{\Omega} g \Delta f \mathrm{~d} x=\sum_{i h \in \Omega_{h}} g(i h)(\Delta f)(i h) h^{3}+\mathrm{O}\left(h^{2}\right) \tag{9}
\end{equation*}
$$

by comparing the integral over the cell centered at $i h$ with the term in the sum. If the cell intersects $\Gamma$, the error in the integrand is $\mathrm{O}(h)$, since g is continuous and has bounded derivative. There are $\mathrm{O}\left(h^{-2}\right)$ such cells, contributing a total error of $\mathrm{O}\left(h \cdot h^{3} \cdot h^{-2}\right)=\mathrm{O}\left(h^{2}\right)$. On each remaining cell the error in the integral is $\mathrm{O}\left(h^{2} \cdot h^{3}\right)$, since $g$ and $\Delta f$ are $C^{2}$. The total error for these cells is $\mathrm{O}\left(h^{2} \cdot h^{3} \cdot h^{-3}\right)=\mathrm{O}\left(h^{2}\right)$ and the claim (9) is verified.

We now have

$$
\begin{equation*}
\int_{\Gamma} f \mathrm{~d} S=\sum_{\Omega_{h}} g \Delta f h^{3}+\mathrm{O}\left(h^{2}\right)=\sum_{\Omega_{h}} g \Delta_{h} f h^{3}+\mathrm{O}\left(h^{2}\right) \tag{10}
\end{equation*}
$$

since $\Delta_{h} f=\Delta f+\mathrm{O}\left(h^{2}\right)$. We can sum by parts and use (7) to obtain

$$
\begin{equation*}
\sum_{\Omega_{h}} g \Delta_{h} f h^{3}=\sum_{\Omega_{h}}\left(\Delta_{h} g\right) f h^{3}=\sum_{\Omega_{h}}\left(w^{h}+\mathrm{O}_{\Gamma}(h)+\mathrm{O}\left(h^{2}\right)\right) f h^{3} . \tag{11}
\end{equation*}
$$

The $\mathrm{O}_{\Gamma}(h)$ error contributes a term of order $h \cdot h^{3} \cdot h^{-2}=h^{2}$ and thus is negligible, as is the other error inside. Combining (10) and (11), we arrive at the conclusion (2).

The fact that the integral is accurate to $\mathrm{O}\left(h^{2}\right)$ although $\tau^{h}=\mathrm{O}(h)$ on the irregular points is related to a gain in accuracy that has long been noted for solutions of elliptic problems using the methods of [3-5,7,8,13]. Proofs of this phenomenon have been given in $[1,4,9]$ and elsewhere. Closely related to the Green's function $g$ solving (4) is the discrete version $g^{h}$ which solves

$$
\begin{equation*}
\Delta_{h} g^{h}=w^{h} \quad \text { in } \Omega_{h}, \quad g=0 \quad \text { on } \partial \Omega_{h} . \tag{12}
\end{equation*}
$$

In fact $g^{h}-g=\mathbf{O}\left(h^{2}\right)$ uniformly; this follows from analytical results in $[1,9]$.

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