

Available online at www.sciencedirect.com



JOURNAL OF COMPUTATIONAL PHYSICS

Journal of Computational Physics 227 (2008) 2195-2197

www.elsevier.com/locate/jcp

Short Note

## A proof that a discrete delta function is second-order accurate

J. Thomas Beale

Department of Mathematics, Duke University, Box 90320, Durham, NC 27708-0320, United States

Received 8 August 2007; accepted 5 November 2007 Available online 17 November 2007

## Abstract

It is proved that a discrete delta function introduced by Smereka [P. Smereka, The numerical approximation of a delta function with application to level set methods, J. Comput. Phys. 211 (2006) 77–90] gives a second-order accurate quadrature rule for surface integrals using values on a regular background grid. The delta function is found using a technique of Mayo [A. Mayo, The fast solution of Poisson's and the biharmonic equations on irregular regions, SIAM J. Numer. Anal. 21 (1984) 285–299]. It can be expressed naturally using a level set function. © 2007 Elsevier Inc. All rights reserved.

© 2007 Elsevier me. 7 m rights reserved.

Keywords: Discrete delta function; Level set function; Surface integral

There is considerable interest in designing accurate discrete delta functions for surfaces in a domain covered by a rectangular grid. They can provide quadrature rules for surface integrals using values at regular grid points [2,10–12]. Such a rule is especially useful when the surface is represented by a level set function. In [10] Smereka constructed a discrete delta function as the truncation error in applying the discrete Laplacian to a "Green's function" for the exact delta function on the surface. To find the truncation error, he used the technique of Mayo [7,8] for solving differential equations with interfacial conditions, in which jump conditions are built into the difference operators on a regular grid. (The immersed interface method [3,5], the EJIIM [13,9] and the ghost fluid method [6] are related to Mayo's technique.) Smereka also showed how to express this delta function in terms of a level set function. He conjectured that the resulting quadrature rule for surface integrals is second-order accurate and verified the accuracy in numerical examples. In this note we give a simple proof of this fact.

Suppose  $\Gamma$  is a closed curve in  $\mathbb{R}^2$  or a closed surface in  $\mathbb{R}^3$ , bounding a set which is contained in a rectangular domain  $\Omega$ . The problem is to design a weight function  $w^h$  at grid points on a square grid  $\Omega_h$ , concentrated near  $\Gamma$ , so that, for any smooth function f defined near the curve  $\Gamma$  in  $\mathbb{R}^2$ ,

DOI of original article: 10.1016/j.jcp.2005.05.005.

E-mail address: beale@math.duke.edu

URL: http://math.duke.edu/faculty/beale

<sup>0021-9991/\$ -</sup> see front matter @ 2007 Elsevier Inc. All rights reserved. doi:10.1016/j.jcp.2007.11.004

J.T. Beale / Journal of Computational Physics 227 (2008) 2195-2197

$$\int_{\Gamma} f(x) \, \mathrm{d}s(x) = \sum_{ih \in \Omega_h} f(ih) w^h(ih) h^2 + \mathcal{O}(h^2) \tag{1}$$

or near the surface  $\Gamma$  in  $\mathbb{R}^3$ ,

$$\int_{\Gamma} f(x) \,\mathrm{d}S(x) = \sum_{ih\in\Omega_h} f(ih)w^h(ih)h^3 + \mathrm{O}(h^2). \tag{2}$$

Arclength and surface area are special cases. Smereka's  $w^h$  has support on the grid points *ih* within distance *h* of  $\Gamma$ , i.e.,  $w^h(ih) = 0$  at other points. We will prove that (2) holds, with  $w^h$  as in [10], assuming  $\Gamma$  is a smooth surface in  $\mathbb{R}^3$ . The case of a curve in  $\mathbb{R}^2$  is entirely similar.

Smereka's procedure is as follows: Let  $\delta_{\Gamma}$  be the distribution, or generalized function, restricting to  $\Gamma$ ; that is, for smooth f on  $\Omega$ ,

$$\int_{\Omega} f \,\delta_{\Gamma} \,\mathrm{d}x = \int_{\Gamma} f \,\mathrm{d}S. \tag{3}$$

Let g be the solution of

$$\Delta g = \delta_{\Gamma} \quad \text{in } \Omega, \quad g = 0 \quad \text{on} \quad \partial\Omega. \tag{4}$$

Assuming  $\Gamma$  is smooth, g is piecewise smooth, i.e., smooth and harmonic on each region bounded by  $\Gamma$ , with the jump conditions

$$[g] = 0, \quad [\partial_n g] = 1 \text{ on } \Gamma, \tag{5}$$

where  $\partial_n$  is the normal derivative on  $\Gamma$ . In fact g can be thought of as a single layer potential on  $\Gamma$ . Now let  $\Delta_h$  be the usual second-order discrete Laplacian on  $\Omega_h$  and let  $\tau^h$  be the truncation error

$$\Delta_h g = \tau^h \quad \text{on } \Omega_h. \tag{6}$$

Smereka constructs the weights  $w^h$  from expressions for  $\tau^h$ , using Mayo's technique [7,8]. At a *regular* grid point  $ih \in \Omega_h$ , for which the stencil of  $\Delta_h$  does not cross  $\Gamma$ ,  $\tau^h(ih) = O(h^2)$  as usual. At an *irregular* grid point,  $\tau^h$  is larger. It can be found to O(h) using the jumps in first and second derivatives of g; see (30) in [10]. These in turn can be expressed in derivatives of the normal and tangent vectors to  $\Gamma$ . (See (41), (47) in [10] for  $\mathbb{R}^2$  and Section 7.2 for  $\mathbb{R}^3$ .) Thus  $\tau^h$  has the form

$$\Delta_h g = \tau^h = w^h + \mathcal{O}_\Gamma(h) + \mathcal{O}(h^2) \quad \text{on } \Omega_h, \tag{7}$$

where  $w^h$  is known analytically and  $w^h$  and  $O_{\Gamma}(h)$  are nonzero only at the irregular points. The errors are uniform. Smereka shows how to write  $w^h$  in terms of a level set function; see (45) and Section 7 in [10].

To prove that (2) is valid, we may assume f is nonzero only in a neighborhood of  $\Gamma$ , as well as smooth. We begin by writing

$$\int_{\Gamma} f \, \mathrm{d}S = \int_{\Omega} f \, \delta_{\Gamma} \, \mathrm{d}x = \int_{\Omega} f \, \Delta g \, \mathrm{d}x = \int_{\Omega} g \Delta f \, \mathrm{d}x. \tag{8}$$

(This could be rewritten in an equivalent way using the jump conditions (5) rather than  $\delta_{\Gamma}$ ).

Next we replace the last integral by a sum over grid points. We check that

$$\int_{\Omega} g\Delta f \, \mathrm{d}x = \sum_{ih\in\Omega_h} g(ih)(\Delta f)(ih)h^3 + O(h^2) \tag{9}$$

by comparing the integral over the cell centered at *ih* with the term in the sum. If the cell intersects  $\Gamma$ , the error in the integrand is O(h), since g is continuous and has bounded derivative. There are  $O(h^{-2})$  such cells, contributing a total error of  $O(h \cdot h^3 \cdot h^{-2}) = O(h^2)$ . On each remaining cell the error in the integral is  $O(h^2 \cdot h^3)$ , since g and  $\Delta f$  are  $C^2$ . The total error for these cells is  $O(h^2 \cdot h^3 \cdot h^{-3}) = O(h^2)$  and the claim (9) is verified.

We now have

$$\int_{\Gamma} f \, \mathrm{d}S = \sum_{\Omega_h} g \Delta f h^3 + \mathcal{O}(h^2) = \sum_{\Omega_h} g \Delta_h f h^3 + \mathcal{O}(h^2) \tag{10}$$

2196

since  $\Delta_h f = \Delta f + O(h^2)$ . We can sum by parts and use (7) to obtain

$$\sum_{\Omega_h} g \Delta_h f h^3 = \sum_{\Omega_h} (\Delta_h g) f h^3 = \sum_{\Omega_h} \left( w^h + \mathcal{O}_{\Gamma}(h) + \mathcal{O}(h^2) \right) f h^3.$$
(11)

The  $O_{\Gamma}(h)$  error contributes a term of order  $h \cdot h^3 \cdot h^{-2} = h^2$  and thus is negligible, as is the other error inside. Combining (10) and (11), we arrive at the conclusion (2).

The fact that the integral is accurate to  $O(h^2)$  although  $\tau^h = O(h)$  on the irregular points is related to a gain in accuracy that has long been noted for solutions of elliptic problems using the methods of [3–5,7,8,13]. Proofs of this phenomenon have been given in [1,4,9] and elsewhere. Closely related to the Green's function g solving (4) is the discrete version  $g^h$  which solves

$$\Delta_h g^h = w^h \quad \text{in } \Omega_h, \quad g = 0 \quad \text{on } \partial\Omega_h. \tag{12}$$

In fact  $g^h - g = O(h^2)$  uniformly; this follows from analytical results in [1,9].

## Acknowledgment

This material is based upon work supported by the National Science Foundation under grant DMS-0404765.

## References

- J.T. Beale, A. Layton, On the accuracy of finite difference methods for elliptic problems with interfaces, Commun. Appl. Math. Comput. Sci. 1 (2006) 91–119. http://www.camcos.org.
- [2] B. Engquist, A.-K. Tornberg, R. Tsai, Discretization of Dirac delta functions in level set methods, J. Comput. Phys. 207 (2005) 28-51.
- [3] R.J. LeVeque, Z. Li, The immersed interface method for elliptic equations with discontinuous coefficients and singular sources, SIAM J. Numer. Anal. 31 (1994) 1019–1044.
- [4] Z. Li, K. Ito, Maximum principle preserving schemes for interface problems with discontinuous coefficients, SIAM J. Sci. Comput. 23 (2001) 339–361.
- [5] Z. Li, K. Ito, The Immersed Interface Method, SIAM, Philadelphia, 2006.
- [6] X.-D. Liu, R. Fedkiw, M. Kang, A boundary condition capturing method for Poisson's equation on irregular domains, J. Comput. Phys. 160 (2000) 151–178.
- [7] A. Mayo, The fast solution of Poisson's and the biharmonic equations on irregular regions, SIAM J. Numer. Anal. 21 (1984) 285-299.
- [8] A. Mayo, The rapid evaluation of volume integrals of potential theory on general regions, J. Comput. Phys. 100 (1992) 236-245.
- [9] V. Rutka, Immersed Interface Methods for Elliptic Boundary Value Problems, Dissertation, T.U. Kaiserslautern, 2005.
- [10] P. Smereka, The numerical approximation of a delta function with application to level set methods, J. Comput. Phys. 211 (2006) 77– 90.
- [11] A.-K. Tornberg, B. Engquist, Numerical approximations of singular source terms in differential equations, J. Comput. Phys. 200 (2004) 462–488.
- [12] J. Towers, Two methods for discretizing a delta function supported on a level set, J. Comput. Phys. 220 (2007) 915-931.
- [13] A. Wiegmann, K.P. Bube, The explicit-jump immersed interface method: finite difference methods for PDEs with piecewise smooth solutions, SIAM J. Numer. Anal. 37 (2000) 827–862.